

Primordial Black Holes from Supercooled Phase Transitions with Radiative Symmetry Breaking

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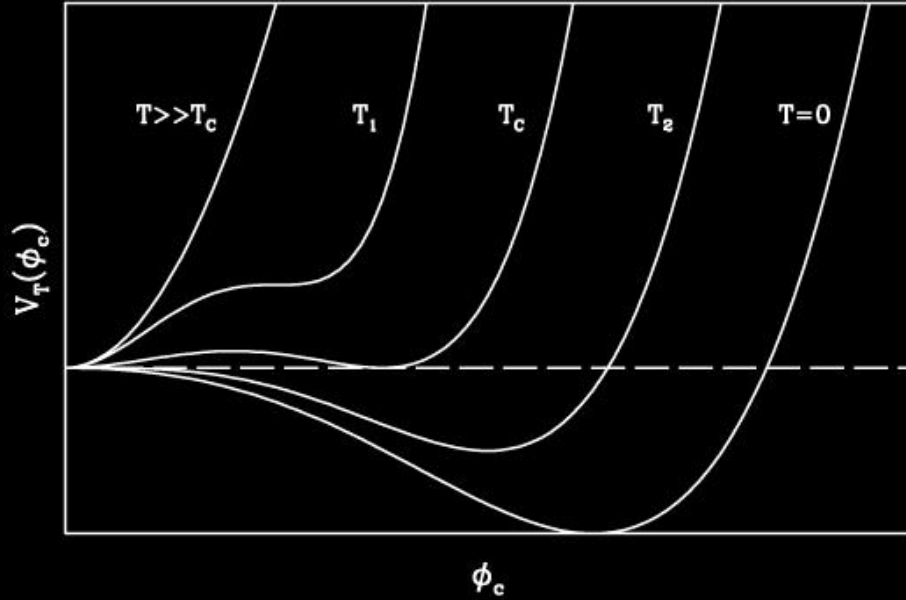
23/01/2025

From Big Bang to Now: a Theory-Experiment Dialogue
SRM University- Andhra Pradesh



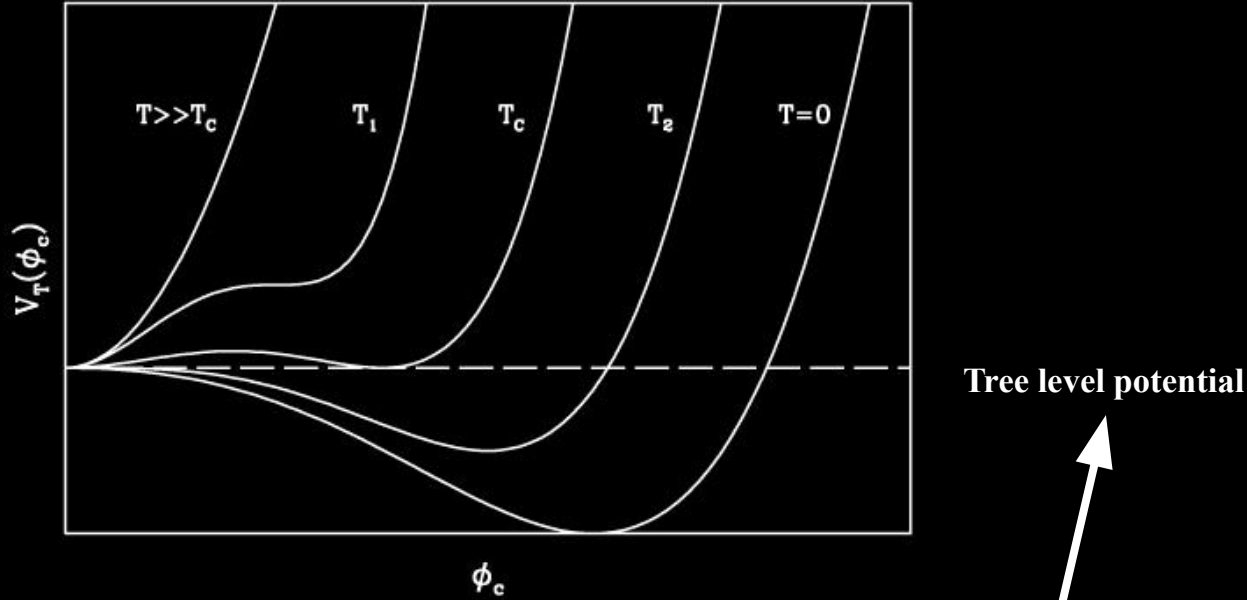
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MINISTRY OF
EDUCATION

First Order Phase Transitions



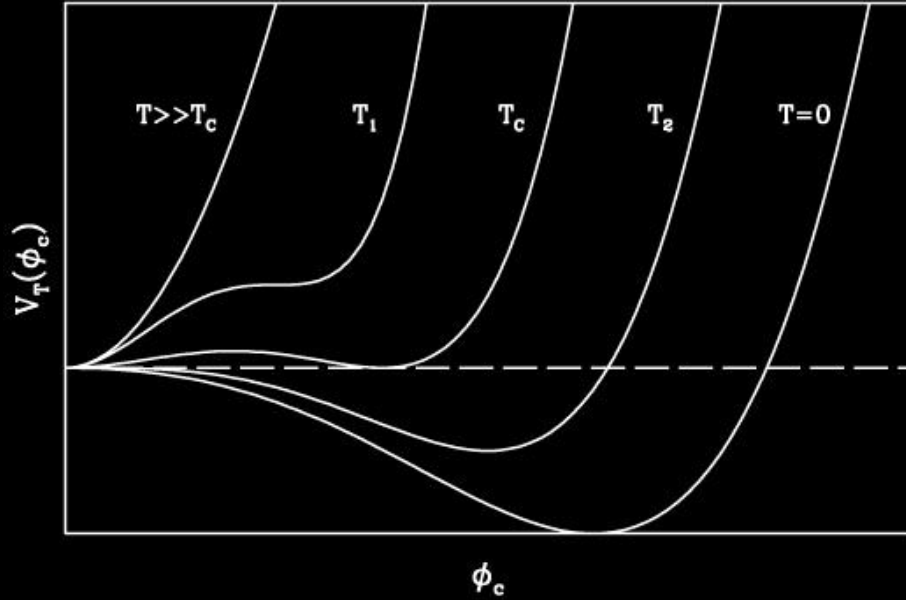
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First Order Phase Transitions



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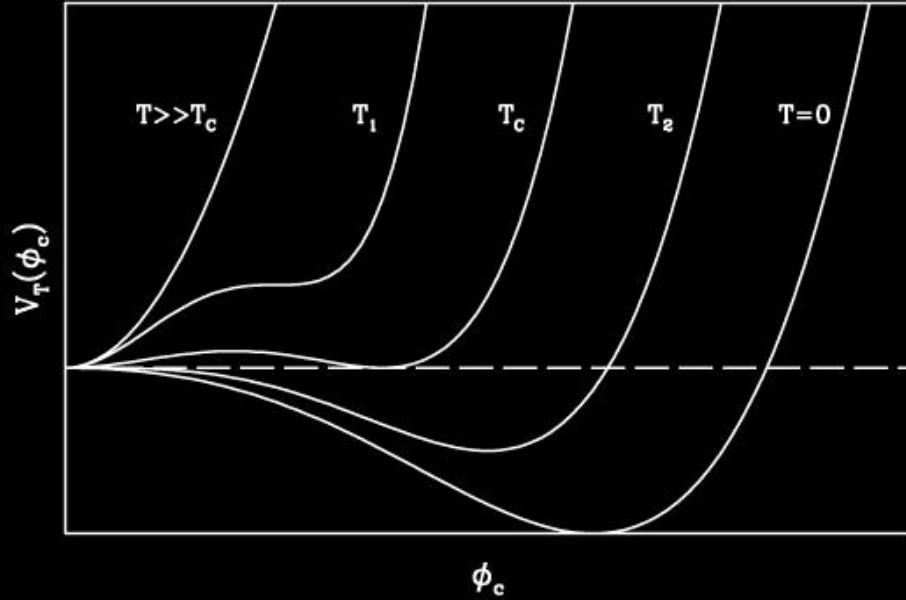


Tree level potential

1-loop corrections
(radiative)

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Finite Temperature
Corrections

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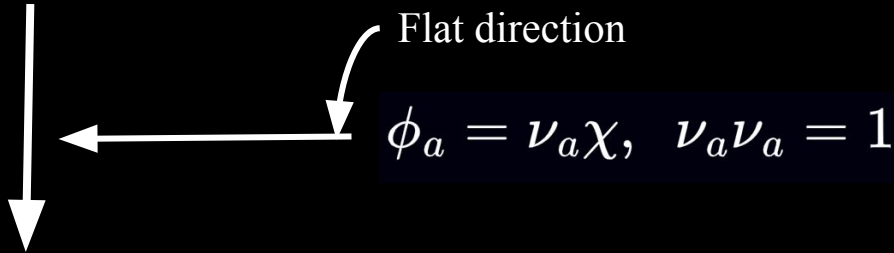
FOPT from Radiative Symmetry Breaking

$$\frac{\lambda_{abcd}}{4!} \phi_a \phi_b \phi_c \phi_d \longrightarrow \text{Tree level potential}$$

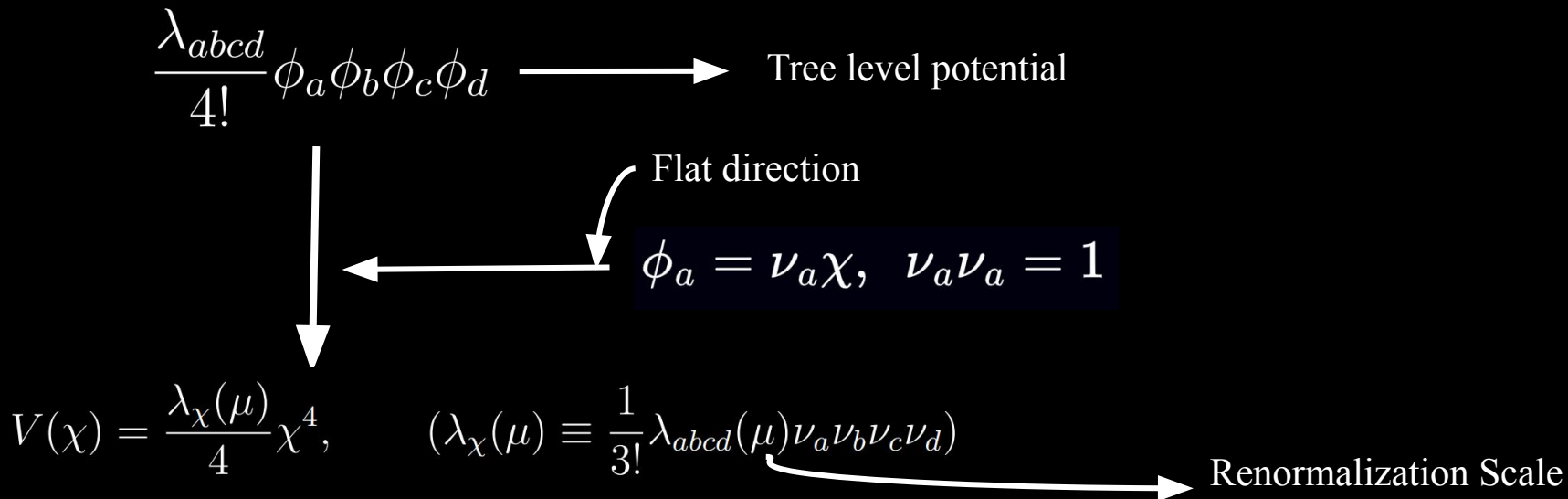
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$$\phi_a = \nu_a \chi, \quad \nu_a \nu_a = 1$$


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$$V(\chi) = \frac{\lambda_\chi(\mu)}{4} \chi^4, \quad (\lambda_\chi(\mu) \equiv \frac{1}{3!} \lambda_{abcd}(\mu) \nu_a \nu_b \nu_c \nu_d)$$

One-loop correction

Renormalization Scale

$$V_q(\chi) = \frac{\bar{\beta}}{4} \left(\log \frac{\chi}{\chi_0} - \frac{1}{4} \right) \chi^4 \quad \text{with} \quad \bar{\beta} \equiv \left[\mu \frac{d\lambda_\chi}{d\mu} \right]_{\mu=\tilde{\mu}} \quad \text{and} \quad \lambda_\chi(\tilde{\mu}) = 0$$

FOPT from Radiative Symmetry Breaking

Including finite temperature corrections, under supercooled expansion

$$\bar{V}_{\text{eff}}(\chi, T) = \frac{m^2(T)}{2}\chi^2 - \frac{k(T)}{3}\chi^3 - \frac{\lambda(T)}{4}\chi^4 + \dots$$

where

$$m^2(T) \equiv \frac{g^2 T^2}{12}, \quad k(T) \equiv \frac{\tilde{g}^3 T}{4\pi}, \quad \lambda(T) \equiv \bar{\beta} \log \frac{\chi_0}{T},$$

$$g^2 \equiv \sum_b n_b m_b^2(\chi)/\chi^2 + \sum_f m_f^2(\chi)/\chi^2, \quad \tilde{g}^3 \equiv \sum_b n_b m_b^3(\chi)/\chi^3$$

also $\epsilon \equiv \frac{g^4}{6\bar{\beta} \log \frac{\chi_0}{T}} \longrightarrow \text{measure of the supercooled expansion}$

FOPT : Important Parameters

- Nucleation Temperature:

$$N(T_n) = \int_{T_n}^{T_c} \frac{dT}{T} \frac{\Gamma(T)}{H(T)^4} = 1 \quad \text{and} \quad \Gamma(T) = \left(\frac{S_3}{2\pi T} \right)^{3/2} T^4 e^{-S_3/T}$$

- Strength of the FOPT:

- Inverse duration:

- Reheating Temperature:

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$$T_{\text{eq}} = T_n \alpha^{1/4}$$

FOPT : Important Parameters

- Nucleation Temperature: $T_n \approx \chi_0 \exp\left(\frac{\sqrt{c'^2 - 16a} - c'}{8}\right)$ $a \equiv \frac{c_3 g}{\sqrt{12}\bar{\beta}}$, $c' \equiv 4 \log \frac{4\sqrt{3}\bar{M}_P}{\sqrt{\bar{\beta}} \chi_0}$

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$$S_3 \approx 27\pi m^3 \frac{1 + \exp(-k/(m\sqrt{\lambda}))}{2k^2 + 9\lambda m^2}$$

- Strength of the FOPT:

$$\alpha = \left(\frac{\Delta V}{\rho_r}\right)_{T=T_n} \longrightarrow \alpha \gg 1 \quad \text{Falls out of calculation}$$

- Inverse duration:

$$\beta/H = T_n \frac{d}{dT} (S_3/T)|_{T=T_n} \longrightarrow \frac{\beta}{H_n} \approx \frac{a}{\log^2(\chi_0/T_n)} - 4$$

- Reheating Temperature:

$$T_{\text{eq}} = T_n \alpha^{1/4} \longrightarrow T_{\text{eq}}^4 = \frac{15\bar{\beta}\chi_0^4}{8\pi^2 g_*(T_{\text{eq}})}$$

First Order Phase Transitions

Significant Observables ?

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Talk by Prof. Sriramkumar

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- **Primordial Black Holes** : collapse of overdense regions due to delayed vacuum decay¹, collapse of curvature perturbations created from time fluctuations of bubble nucleation²

¹2106.05637, 2212.14037, 2305.04924, 2305.04942, ...

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Main point of this talk!

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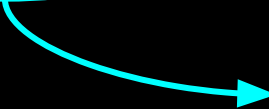
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Detour: Why PBH?

- PBHs can partially or completely play the role of DM.

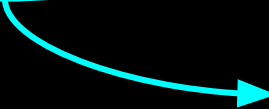
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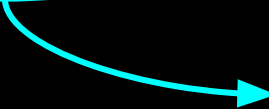
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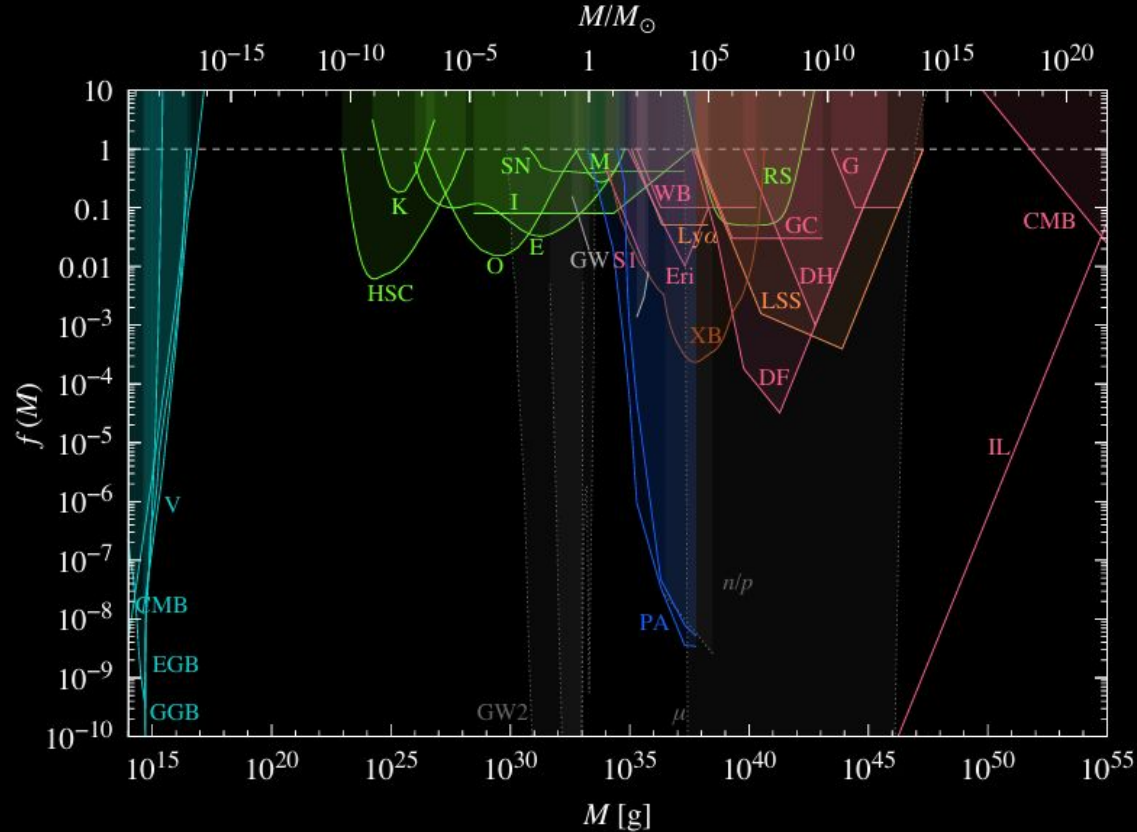
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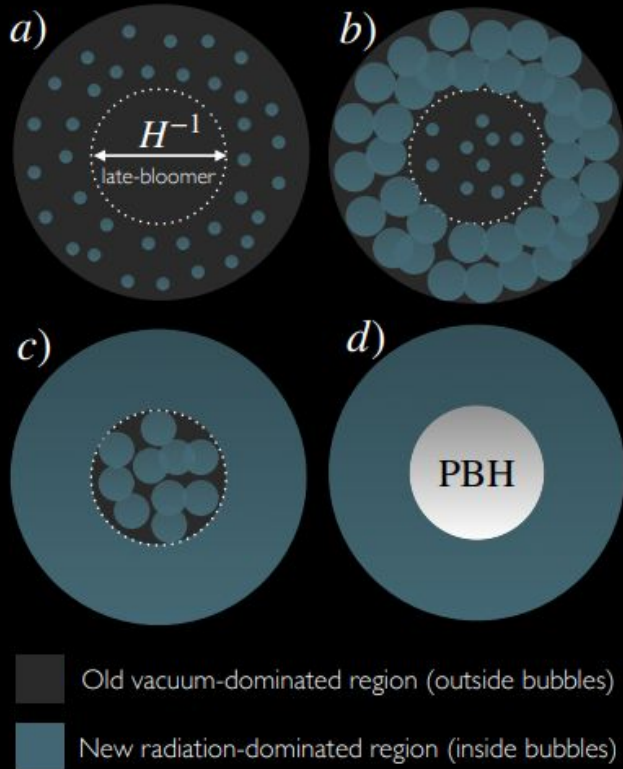
- PBH properties can help us gain insight regarding very early universe
- PBH might also be the only way for us to detect Hawking evaporation

Detour: Why PBH?



Talk by Prof. Sriramkumar

FOPT: PBH Creation



- Delayed vacuum decay in some region
- Background energy density gets diluted with the expansion of the universe
- Vacuum energy remains constant and eventually creates a overdense region
- Depending on the overdensity, it may collapse into a PBH

FOPT: PBH Creation

- Collapse Probability $\mathcal{P}_{\text{coll}} \approx \exp \left[-a_{\mathcal{P}} \left(\frac{\beta}{H_n} \right)^{b_{\mathcal{P}}} (1 + \delta_c)^{c_{\mathcal{P}} \frac{\beta}{H_n}} \right]$
 $a_{\mathcal{P}} \approx 0.5646, b_{\mathcal{P}} \approx 1.266, c_{\mathcal{P}} \approx 0.6639$

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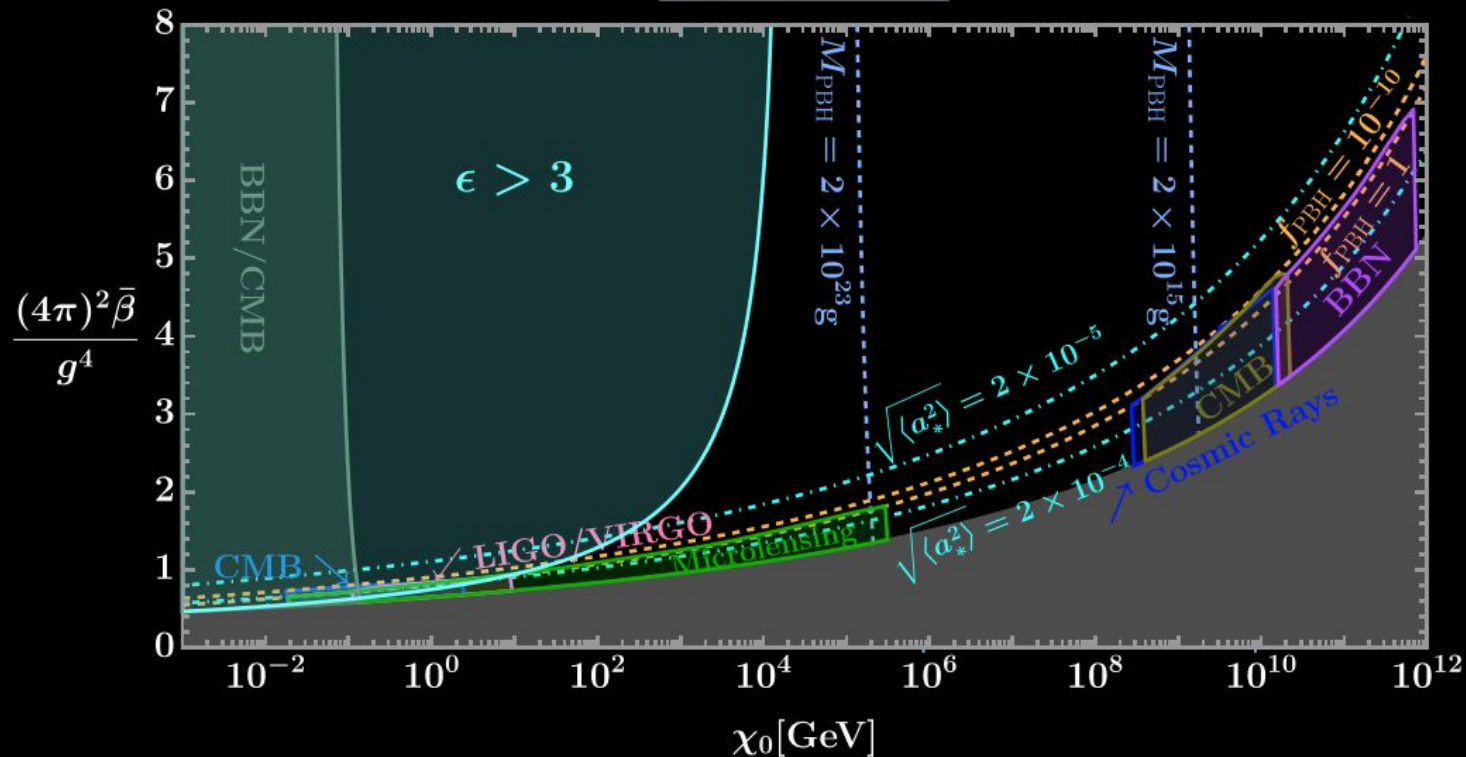
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- Initial PBH spin

$$\sqrt{\langle a_*^2 \rangle} \approx \frac{2.1 \times 10^{-3}}{23.484 - 1.25 \log_{10}(f_{\text{PBH}}) - 1.25 \log_{10} \left(\frac{\Omega_{\text{CDM}}}{0.26} \right) - 0.625 \log_{10} \left(\frac{M_{\text{PBH}}}{10^{15} \text{ g}} \right)}$$

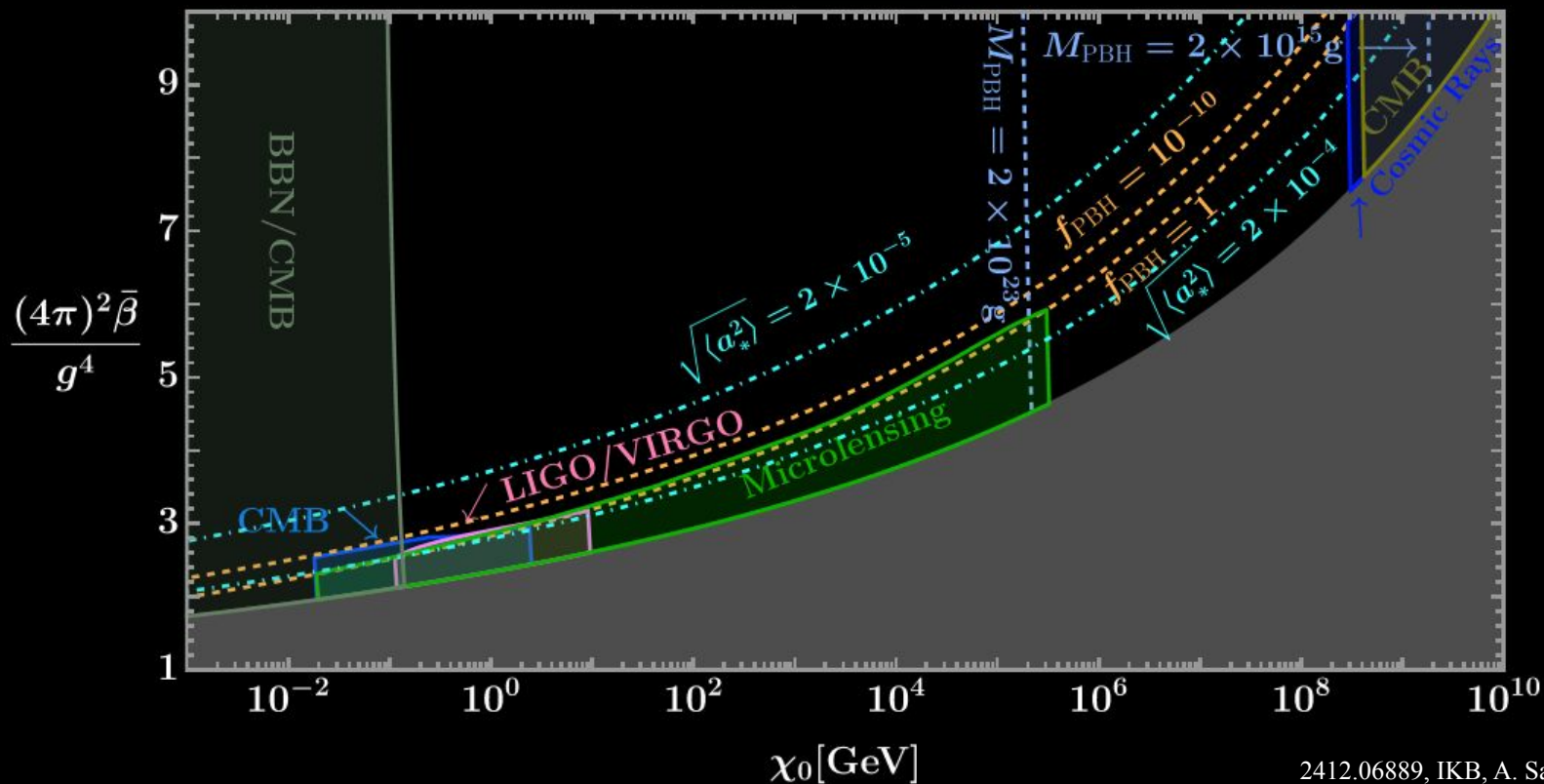
Results

$$g = \tilde{g} = 0.7$$



Results

$$g = \tilde{g} = 0.5$$



Summary and Future Prospects

- RSB scenario is well motivated by various BSM extensions
- It is also the most likely scenario through which an FOPT can create PBHs
- However, this high supercooling can lead to small period of non-standard cosmology with interesting implications.
- The spin of the PBHs can also be impacted by non-standard cosmological scenario.

THANK YOU

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One-loop correction

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$$\mu \frac{dV_q}{d\mu} = 0, \quad \text{where} \quad V_q \equiv V + V_1 + V_2 + \dots$$

A specific example: $U(1)_{B-L}$

$$\begin{aligned} \mathcal{L}_{\text{SM}}^{\text{ns}} + D_\mu A^\dagger D^\mu A + \bar{N}_j i \not{D} N_j - \frac{1}{4} B'_{\mu\nu} B'^{\mu\nu} \\ + \left(Y_{ij} L_i \mathcal{H} N_j + \frac{1}{2} y_{ij} A N_i N_j + \text{h.c.} \right) - \lambda_a |A|^4 + \lambda_{ah} |A|^2 |\mathcal{H}|^2 \end{aligned}$$

$$D_\mu = \partial_\mu + i g_3 T^\alpha G_\mu^\alpha + i g_2 T^a W_\mu^a + i g_Y \mathcal{Y} B_\mu + i [g_m \mathcal{Y} + g'_1 (B - L)] B'_\mu$$

$$(4\pi)^2 \mu \frac{d}{d\mu} \lambda_a = 96 g_1'^4 - 48 \lambda_a g_1'^2 + 20 \lambda_a^2 + 2 \lambda_{ah}^2 + 2 \lambda_a \text{Tr}(y y^\dagger) - \text{Tr}(y y^\dagger y y^\dagger)$$

$$\bar{\beta} = \frac{96 g_1'^4}{(4\pi)^2} \qquad g = 2\sqrt{3} |g'_1|$$